

## Homework 3: Orthogonal Coordinate Systems, Velocity and Acceleration

Due Monday, February 5

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### Problem 1: Velocity and acceleration in SPC

Using your results from the previous homework, derive expressions for the velocity ( $\dot{\vec{r}}$ ) and acceleration ( $\ddot{\vec{r}}$ ) vectors in spherical polar coordinates.

### Problem 2: Line integrals in polar coordinates

Express the vector  $\vec{A} = x^2y \hat{x} + y^2x \hat{y}$  entirely in polar coordinates (including the unit vectors!), and then evaluate the integral of  $\vec{A} \cdot d\vec{\ell}$  along the following paths:

- (a) From  $\phi = 0$  to  $\phi = \pi/2$  along the path  $\rho = 1$ .
- (b) From  $\rho = 1$  to  $\rho = 3$  along the path  $\phi = \pi/4$ .

### Problem 3: Elliptical cylindrical coordinates

Elliptical cylindrical coordinates describe points using a grid composed of ellipses and hyperbolae in the  $x$ - $y$  plane, along with the usual  $z$  coordinate. They are related to Cartesian coordinates by

$$x = a \cosh \mu \cos \nu \quad (1)$$

$$y = a \sinh \mu \sin \nu \quad (2)$$

$$z = z, \quad (3)$$

where  $a$  is a positive number,  $0 \leq \mu < \infty$ , and  $0 \leq \nu < 2\pi$ . Derive the scale factors and unit vectors for these coordinates.

**HINT:** If you haven't seen  $\cosh$  and  $\sinh$  in a while, they are defined as

$$\cosh \mu = \frac{1}{2} (e^\mu + e^{-\mu}) \quad \sinh \mu = \frac{1}{2} (e^\mu - e^{-\mu}) . \quad (4)$$

Since  $d(e^{\pm\mu})/d\mu = \pm e^{\pm\mu}$ , their derivatives are

$$\frac{d \cosh \mu}{d\mu} = \sinh \mu \quad \frac{d \sinh \mu}{d\mu} = \cosh \mu . \quad (5)$$

You can also check that  $\cosh^2 \mu - \sinh^2 \mu = 1$  by plugging in the expressions above and expanding. So  $\cosh$  and  $\sinh$  are like  $\cos$  and  $\sin$ , with plusses where you'd expect minuses and vice-versa.

### Problem 4: Celestial mechanics

In class we looked at Newton's laws applied to the problem of a planet orbiting a star. To a pretty good approximation, the planet's distance  $r(t)$  from the star satisfies the equation

$$\ddot{r}(t) - \frac{j^2}{r(t)^3} = -\frac{MG}{r(t)^2}, \quad (6)$$

where  $j$  is the planet's angular momentum divided by its mass,  $M$  is the mass of the star, and  $G$  is Newton's constant. The planet's angular momentum  $J = m r(t)^2 \dot{\phi}(t)$  stays the same throughout

the orbit, so  $j = J/m = r(t)^2 \dot{\phi}(t)$  is constant. We saw in class that we could rewrite this as an equation for  $r(\phi)$  instead of  $r(t)$ . Using the chain rule and the relationship between  $\dot{\phi}$ ,  $r$ , and  $j$ , the time derivative of  $r$  can be written as

$$\dot{r}(t) = \frac{dr}{dt} = \frac{d\phi}{dt} \frac{dr}{d\phi} = \frac{j}{r^2} \frac{dr}{d\phi}. \quad (7)$$

Show that once we rewrite equation (6) this way<sup>1</sup> and apply the substitution  $r(\phi) = 1/u(\phi)$ , we wind up with the much simpler equation

$$u''(\phi) + u(\phi) = \frac{MG}{j^2}, \quad (8)$$

where  $'$  indicates a derivative with respect to  $\phi$ .

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<sup>1</sup>You'll have to apply the chain rule twice for the  $\ddot{r}(t)$  term!